

STRUCTURAL, ELECTRONIC AND OPTICAL PROPERTIES OF CUBIC PEROVSKITE CsPbX₃ (X= Br, Cl And I)

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ABSTRACT:

Plane waves with norm conserving pseudopotentials (PW-PP) method in conjunction with density functional theory (DFT) frame work have been used to investigate structural, electronic and optical properties of lead-halide cubic perovskite CsPbX₃ (X=Br, Cl and I). The generalized gradient approximation (GGA), specifically Perdew-Burke-Ernzerhof (PBE) flavor, has been chosen to treat the exchange correlation term of Kohn-Sham equation. Structural parameters are comparable with other theoretical and experimental studies. In spite of good agreement of our band gap values E_g with other theoretical works, however, they were not comparable when compared to the experimental E_g values due to the well-known problem of E_g value underestimation of DFT. To update the E_g value, we have used GW method as a self-consistent quasiparticle method on energies and wave functions and indeed they have been improved. Optical properties have been calculated using density functional perturbation theory (DFPT). Our results show that CsPbX₃ (X=Br, Cl, I) has maximum response to the electromagnetic spectrum at low energies (visible region) but minimum response at high energies.

KEYWORDS: Ab initio calculation; Structural, electrical and optical properties; Band gap; Perovskites CsPbX₃.

1. INTRODUCTION

Perovskite-based compounds (with formula ABX₃) have been studied intensively in the recent years as promising materials for light emitting diodes (LED) (Song et al., 2015; Tan et al., 2014), photodetectors (Dong et al., 2015; Wei et al., 2016) and for low-cost solar cells with relatively high efficiency, up to 22.1% (Kim et al., 2012; Kojima, Teshima, Shirai, & Miyasaka, 2009; Y. Li, Duan, Zhao, & Tang, 2018; Mutalikdesai & Ramasesha, 2017; Zhou et al., 2014), because of their good electronic and optical properties. However, potential stability is the most challenge which hamper mass production of perovskite-based solar cells (Z. Li et al., 2015). Organometallic halide perovskite (OHP) (where A is organic cation, B is metal cation and X is halide anion) and inorganic halide perovskite (IHP) (for example, A=Cs) have been studied. In spite of instability of both groups, the latter group is more promising because the OHP are highly sensitive to moisture and other ambient environment conditions (X. Li et al., 2017). Among IHP, CsPbX₃ (X=Br, Cl and I) are direct and tunable band gap semiconductors (Castelli, García-Lastra, Thygesen, & Jacobsen, 2014). The three compounds show phase transition, CsPbBr₃ has orthorhombic structure at room temperature and undergoes phase transition to tetragonal at 88 [°C] and to the cubic at 130 [°C] (MØLLER, 1958; Stoumpos et al., 2013), CsPbCl₃ undergoes phase transition from tetragonal structure or monoclinic to cubic perovskite at 47 °C (MØLLER, 1958) and CsPbI₃, which is orthorhombic, becomes perovskite cubic above 316 °C (Surendra Sharma, Weiden, & Weiss, 1992).

The CsPbBr₃ meets sufficient requirements for photoelectronic devices such as x-ray and γ -ray detectors (Stoumpos et al., 2013) and long-term stability solar cells

with efficiencies 6.95% and 5.37 % (Lei et al., 2018). A cubic perovskite CsPbI₃ solar cell has been fabricated with efficiency up to 2.9% with fair stability at room temperature (Eperon et al., 2015).

A number of theoretical works on CsPbX₃ (X=Br, Cl and I) have been published. Lang et al (Lang, Yang, Liu, Xiang, & Gong, 2014) calculated imaginary part of dielectric function and electronic band structure implementing the projector augmented-wave (PAW) pseudopotentials (Kresse & Joubert, 1999) in the framework of density functional theory (DFT) using the Vienna ab-initio simulation package (VASP) (Kresse & Furthmüller, 1996). Chang and Park (Chang, Park, & Matsuishi, 2004) investigated structural and electronic properties of CsPbX₃ theoretically using first-principles pseudopotential total-energy calculations within the local density approximation (LDA). All electrons full potential linearized augmented plane wave (FP-LAPW) method was implemented to calculate structural, electronic, and optical properties of cubic perovskites CsPbX₃ using the wien2k code (Murtaza & Ahmad, 2011). Structural, electronic and optical properties of three compounds CsPbX₃ have been calculated using full potential linearized augmented plane waves method within DFT and for accurate band gap values, modified Becke-Johnson (TB-mBJ) exchange potential has been employed (Ahmad et al., 2017).

The aim of this study is to investigate structural, electronic and optical properties using DFT frame. The augmented plane waves (PW) with norm-conserving, separable, dual-space Gaussian-type pseudopotentials of Goedecker, Teter and Hutter scheme (GTH) (Hartwigsen, Goedecker, & Hutter, 1998) with generalized gradient approximation (GGA) of DFT (Perdew, Burke, & Ernzerhof, 1996) have been used to calculate structural, electric and optical properties of the perovskites CsPbX₃ (X=Br, Cl, I). Perdew-Burke- Ernzerhof (PBE) has been chosen to treat

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exchange-correlation term of DFT Kohn-Sham total energy equation (Kohn & Sham, 1965). To obtain accurate value of the fundamental band gap, we have used GW approximation (Hedin, 1965).

2. COMPUTATIONAL METHOD

In the present work, DFT framework in conjunction with PW and norm-conserving pseudopotentials has been used to calculate structural, electronic and optical properties of CsPbX₃ (X=Br, Cl and I). The GTH separable, dual-space Gaussian-type pseudopotentials (Krack, 2005) have been employed as it treats 5s² 5p⁶ 6s¹ of Cs, 5d¹⁰ 6s² 6p² of Pb, 4s² 4p⁵ of Br, 3s² 3p⁵ of Cl and 5s² 5p⁵ of I as valence states. Our calculations have been done using Abinit package (Gonze et al., 2009).

Cubic perovskite structure (ABX₃) has 5 atoms per unit cell. The five atoms taken into account are located as follows: A at 0.0 0.0 0.0, B at 0.5 0.5 0.5 and X₃ atoms occupy 0.5 0.5 0.0, 0.5 0.0 0.5 and 0.0 0.5 0.5 (figure 1).

First, convergence calculations have been done to result in the cut-off energy of 1088 [eV] and the number of k points of 40 points, which corresponds to the Monkhorst-Pack (Monkhorst & Pack, 1976) k-point mesh of 6×6×6. Then we have optimized lattice constant *a* and atomic coordinates using GGA-PBE method. The optimized *a* and relaxed atomic coordinates were employed to compute band structure and band gap values *E_g*. Moreover, total density of states (TDOS) and partial density of states (PDOS) have been computed. Due to the familiar underestimation of band gap problem of DFT (Perdew, 1986), GW method (Hedin, 1965) as a self-consistent quasiparticle method on energies and wavefunctions was used to update and to get *E_g* values comparable with experimental results. Optical properties have been calculated using density functional perturbation theory (DFPT) (S Sharma & Ambrosch-Draxl, 2004) and Kramers-Kronig relations (Collins et al., 2001).

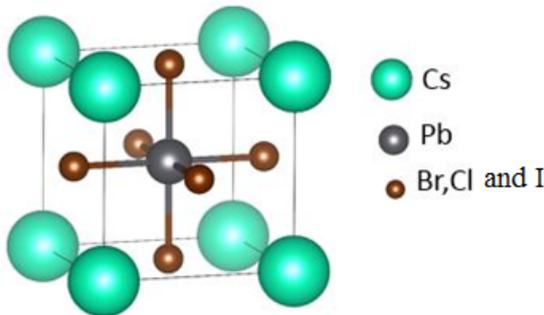


Figure 1. The cubic structure of CsPbX₃ perovskite

3. RESULTS AND DISCUSSION

3.1 Structural properties

Optimization was done for (*a*) and atomic coordinates of the three compounds. Employing the relaxed atomic coordinates, the total energy (*E_T*) and pressure were calculated for various values of (*a*). Figure 2 shows the *E_T* plots against *a*, from which the optimized *a* values for the three compounds were taken, and figure 3 displays the change of pressure with the volume of the unit cell.

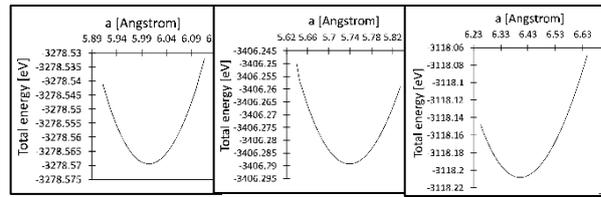


Figure 2. Total energy versus lattice parameter of (a) CsPbBr₃ (b) CsPbCl₃ and (c) CsPbI₃

Figure 2. Total energy versus lattice parameter of (a) CsPbBr₃ (b) CsPbCl₃ and (c) CsPbI₃

The change of pressure *P* with the volume *V* of the unit cell for each compound, see figure 3, was fitted to the second-order Birch-Murnaghan equation of state (Birch, 1947; Murnaghan, 1944):

$$P(V) = \frac{3B_0}{2} \left[\left(\frac{V_0}{V} \right)^{7/3} - \left(\frac{V_0}{V} \right)^{5/3} \right] \left[1 + \frac{3}{4} (B'_0 - 4) \left\{ \left(\frac{V_0}{V} \right)^{2/3} - 1 \right\} \right] \quad (1)$$

Where *V₀* is the cell volume at zero pressure, *B₀* is the bulk modulus and *B'₀* is the derivative of bulk modulus with respect to pressure. From the fittings, we obtained the values of *B₀* of the three compounds and they are given in Table 1 along with their *V₀* values. The optimized structural properties, calculated in GGA-PBE, as well as results of other work are also displayed in Table 1.

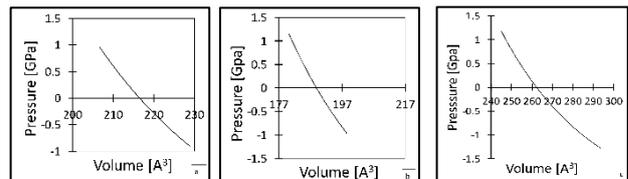


Figure 3. Pressure evolution of unit cell volume of (a) CsPbBr₃, (b) CsPbCl₃ and (c) CsPbI₃.

Table 1. Optimized structural properties calculated in GGA-PBE compared with experimental and other theoretical results

Parameter	CsPbBr ₃			CsPbCl ₃			CsPbI ₃		
	Present work	Exp.	Other theoretical works	Present work	Exp.	Other theoretical works	Present work	Exp.	Other theoretical works
a [Angstrom]	6.007	5.87 ^a	6.046 ^b	5.738	5.6 ^a	5.783 ^b	6.397	6.289 ^c	6.461 ^b
E _T [eV]	-3278.56			-3406.28			-3118.2		
B ₀ [GPa]	18.06		18.1 ^d	22.09		22.59 ^e	13.02		19.8 ^f
V ₀ [Å ³]	216.756			189.012			261.77		
Bonds length [Angstrom]									
Cs-Br	4.245		4.15 ^e						
Pb-Br	3.002		2.935 ^e						
Cs-Cl				4.0578		6.883 ^e			
Pb-Cl				2.8693		2.810 ^e			
Cs-I							4.526		4.137 ^e
Pb-I							3.2		5.406 ^e

a ref.(MØLLER, 1958), b ref. (Ye et al., 2015), c ref. (Trots & Myagkota, 2008), d ref (Zhang, Zeng, & Wang, 2017), e ref. (Hu, Ge, Yu, & Feng, 2017), f ref. (Murtaza & Ahmad, 2011), g ref. (Ahmad et al., 2017).

There is, mostly, good agreement of our results with experimental results and with other theoretical works, which have been done using different implementations of DFT. Lattice parameter values of CsPbBr₃, CsPbCl₃ and CsPbI₃ are in error with experimental values by 2.3%, 2.4% and 1.7%, respectively. The Bulk modulus values of CsPbBr₃, CsPbCl₃ have errors of 0.16% and 2.19%, respectively, while that of CsPbI₃ has an error of 34.24%. The bond lengths are also in very good agreement except for Cs-Cl bond of CsPbCl₃ and Pb-I bond of CsPbI₃.

3.2 Electronic properties

Electronic shape and nature are described by electronic band structure and density of states (DOS), and both have been computed using GGA-PBE method. Figure 4 shows band structure of CsPbBr₃, CsPbCl₃ and CsPbI₃ compounds. The band structures shown in figure 4 confirm that the three compounds are semiconductors. The little difference in the band structure shape of the three compounds is attributed to the difference of halogen atom (Br, Cl and I)

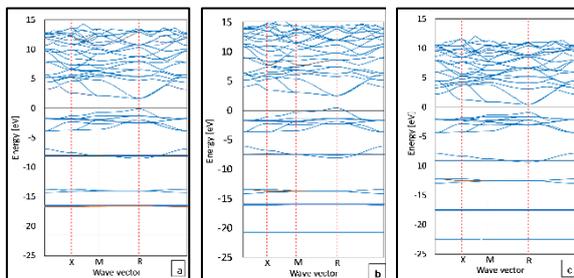


Figure 4. Band structure of (a) CsPbBr₃, (b) CsPbCl₃ and (c) CsPbI₃. The PED-GGA method of DFT was used.

It is confirmed that the cesium based lead-halide perovskites have a direct fundamental band gap E_g between bands 22 and 23. Minimum of the conduction band (CBM) and maximum of the valence band (VBM) both are located at the symmetry point R. Notably, bands 23, 24 and 25 are degenerate, that is, more states for electron transition from VB; this makes these compounds even more significant. Band gap values at high symmetry points including the fundamental band gap are listed in table 2.

Table 2. Calculated energy gaps (in [eV]) at some high symmetry points in PBE-GGA method. Results of other theoretical and experimental works are also given.

	E ^{Γ-Γ}	E ^{X-X}	E ^{M-M}	E ^{R-R}
CsPbBr ₃				
Present work				1.75
Theoretical work	4.91	3.85	2.82	1.6 ^a
Experiment	5.05 ^a	4.5 ^a	3.05 ^a	2.3 ^c
CsPbCl ₃				
Present work	5.69	4.52	3.37	2.18
Theoretical work	5.5 ^a	4.6 ^a	3.2 ^a	1.8 ^a
Experiment				3.0 ^c
CsPbI ₃				
Present work	4.21	3.04	2.29	1.46
Theoretical work	4.2 ^a	3.6 ^a	2.4 ^a	1.3 ^a
Experiment				1.73 ^b

a ref. (Murtaza & Ahmad, 2011), b ref. (Eperon et al., 2014), c ref. (Heidrich et al., 1981).

There is good agreement between our results in Table 2 and other theoretical results, except for the value of the fundamental band gap E_g = E^{R-R}, which does not agree with experimental data. This discrepancy is attributed to the underestimation of the PBE-GGA method of the fundamental band gap value, which is a well-known problem of DFT. GW method has been used to correct the value of E_g. Table 3 shows improvement of the calculated E_g values after implementing the GW method. Fundamental band gap values of CsPbX₃ (X=Br, Cl and I), calculated using PBE-GGA, are in error with experimental values by 25.49%, 27.16%

and 15.4%, respectively, but when we have used GW method, they have been improved, and the errors reduced to 2.54%, 6.66% and 10.8 %, respectively. It is clear from the fundamental band gap values that CsPbCl₃ has largest band gap which can be attributed to the smaller atomic number of Cl compared to those of Br and I atoms. Higher atom number means more number of protons that are electrostatically attracting and hence lowering the bonding energy and band gap as well.

Table 3. Fundamental band gap calculated using PBE-GGA and GW along with results of other theoretical and experimental works.

	$E_g = E^{R-R}$ [eV]		
	PBE-GGA	GW	Experimental
CsPbBr ₃			
Present work	1.75	2.42	2.36 ^c
Other works	1.61 ^a	2.36 ^f	
CsPbCl ₃			
Present work	2.18	3.2	3 ^c
Other works	2.16 ^a	3.27 ^f	
CsPbI ₃			
Present work	1.46	1.91	1.73 ^d
Other works	1.47 ^a , 1.3 ^b	1.96 ^f	

a ref. (Ahmad et al., 2017), b ref. (Murtaza & Ahmad, 2011), c ref. (Kulbak, Cahen, & Hodes, 2015), d ref. (Eperon et al., 2014), e ref. (Gesi, Ozawa, & Hirotsu, 1975), f ref. (Becker et al., 2018).

Total density of states (DOS) and partial density of states (PDOS) are more descriptive of the electronic nature of CsPbX₃ (X=Br, Cl and I) compounds. Figure 5 shows the calculated DOS and PDOS where the latter establishes the contribution of each atom to DOS. The DOS is plotted in the range of -20 to 20 [eV], however the states start to appear at about -17 [eV]. Fermi level in our calculations are located at -0.22 [eV], 0.2 [eV] and -1.11 [eV] for CsPbBr₃, CsPbCl₃ and CsPbI₃, respectively. The Fermi level energy values confirm that the compounds CsPbX₃ (X=Br, Cl and I) are semiconductors.

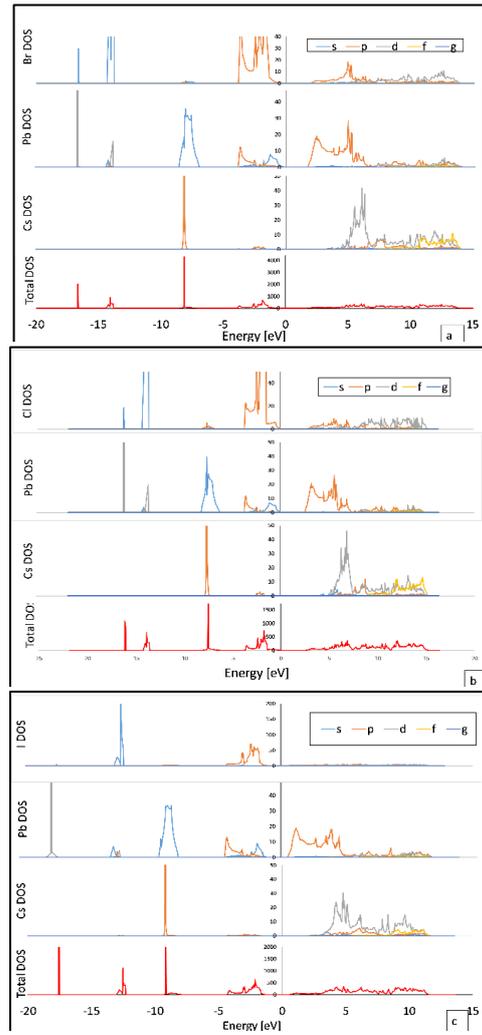


Figure 5. Total and partial DOS of (a) CsPbBr₃, (b) CsPbCl₃ and (c) CsPbI₃ calculated using PBE-GGA functional.

The PDOS describes the contributions of each atomic state to the DOS. For CsPbX₃ (X=Br, Cl and I), DOS at VB is mostly due to the contribution of the p state of the halogen atom (Br, Cl, and I), moreover, p state and s states of Pb are also involved. The CB are mostly due to the contributions of the d state of Cs and p state of Pb, see figure 5.

3-3 Optical properties

The CsPbX₃ (X=Br, Cl and I) are considered promising materials to future efficient solar cells from the point of view of their band gap values (Eperon et al., 2015; Lei et al., 2018; Stoumpos et al., 2013) and high optical response to the electromagnetic spectrum. We have calculated the frequency dependent linear optical properties, namely, imaginary part of dielectric function ($\text{Im}(\epsilon(\omega))$), real part of dielectric function ($\text{Re}(\epsilon(\omega))$), refractive index $n(\omega)$ and absorption coefficient ($\alpha(\omega)$). The results are shown in figure (6).

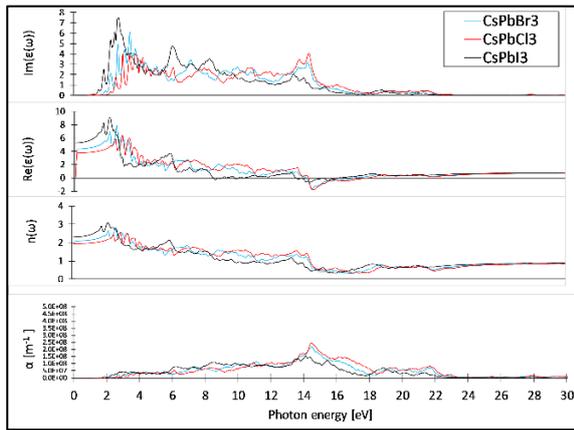


Figure 6. Optical properties of CsPbX₃ (X=Br, Cl and I).

We have investigated optical properties response to the photon energies starting from 0 to 30 [eV] in steps of 0.0186 [eV]. The imaginary part of the dielectric function ($\text{Im}(\epsilon(\omega))$) indicates absorption limits of the materials and energy gain ability as solar cells. The real part ($\text{Re}(\epsilon(\omega))$) indicates energy store ability of the material. Let us divide the ($\text{Im}(\epsilon(\omega))$) curve into two main regions: a low energy region and a high energy region, and discuss it in conjunction with band gap and DOS. In the low energy region, electron transitions occurs from p valence state of the halogen (Br, Cl and I) to the p conduction state of Pb, in this region CsPbI₃ has the highest peak due to its lowest band gap among these materials, therefore, it has maximum response to the visible light. At high energies, the response is due to the electron transitions from mixed s valence state of Pb and p valence state of Cs to the mixed p conduction state of Pb and d and p conduction states of Cs (Murtaza & Ahmad, 2011). In the high energy region, CsPbCl₃ has maximum peak followed by a lower peak for CsPbBr₃ then a lowest for CsPbI₃. This occurs because CsPbCl₃ has higher DOS in the p state of Pb in the VB which contains more electrons, thus leading to more transitions from VB to the CB. Similar trend of explanation can be used for the other optical properties. For example, maximum peak of absorption coefficient $\alpha(\omega)$ at visible light region (low energy) for CsPbCl₃, CsPbBr₃ and CsPbI₃ are 2.59×10^7 , 2.62×10^7 and 4.45×10^7 [m^{-1}], respectively. On the other hand, maximum peaks of $\alpha(\omega)$ in the high energy region are located at 2.42×10^8 , 2.15×10^8 and 1.54×10^8 [m^{-1}] for CsPbCl₃, CsPbBr₃ and CsPbI₃, respectively.

From the real part of dielectric function ($\text{Re}(\epsilon(\omega))$), the dielectric constant values can be found for any frequency. Its values at zero frequency $\epsilon(0)$ are 3.69, 4.32 and 5.56 for CsPbCl₃, CsPbBr₃ and CsPbI₃, respectively. They increase with frequency and reach the maximum values at 2.93, 2.61 and 2.2 [eV], respectively, in the visible light region. Then, they decrease to negative value before they start again to increase slightly into less than one.

Refractive index $n(\omega)$ starts to increase until it reaches the maximum values of 2.54, 2.82 and 3.07 at energies 2.93, 2.61 and 2.2 [eV] for CsPbCl₃, CsPbBr₃ and CsPbI₃, respectively. CsPbI₃ has highest $n(\omega)$ which makes it a better material to retain energy than CsPbBr₃ and CsPbCl₃.

4. CONCLUSIONS

First principle calculations have been done here to investigate and review structural, electronic and optical properties of CsPbX₃ (X=Br, Cl and I). The structural properties have a

good agreement with other theoretical works and experimental results. Moreover, it shows that CsPbX₃ (X=Br, Cl and I) have direct E_g between 22 and 23 bands located at the symmetry point R. The E_g values are comparable to the results obtained experimentally and theoretically using different implementations of DFT method by others.

The bands 23, 24 and 25 have the same energy (degenerate) thus accommodating more places for electrons to be excited from VB, hence leading to higher response to the electromagnetic excitation. Among CsPbX₃ (X=Br, Cl and I) group, CsPbI₃ has the highest response peak at visible light spectrum due to its E_g value. On the other hand, CsPbCl₃ has the highest response peak at high energy spectrum because it has a highest DOS specifically s state of Pb. The latter means more electrons available to transit to the CB.

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